# Substrate Integrated Waveguide Filter With Flat Passband Based on Complex Couplings

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*Abstract***— A lossy filter with flat passband and asymmetric response is proposed by using complex couplings. The complex coupling coefficients with positive and negative imaginary parts are realized with transmission lines loaded with series and shunt resistors, respectively. Design formulas are derived for the critical parameters. A filter prototype is developed at 4.97 GHz and substrate integrated waveguide resonators with unloaded** *Q***-factor of 450 are used. One transmission zero is located in its upper stopband. The 0.2-dB bandwidth of 279 MHz is achieved and the equivalent** *Q***-factor is about 1900. Good agreement is obtained among the synthesized, simulated, and measured results.**

*Index Terms***— Complex coupling, lossy bandpass filter, substrate integrated waveguide (SIW), uniform** *Q***-factor.**

### I. INTRODUCTION

**T**HE design of high-performance miniaturized bandpass fil-<br>ters is always a challenge, due to relatively low unloaded<br> $\Omega$  feature of sommast meanstons. Low  $\Omega$  filters suffer from *Q*-factors of compact resonators. Low-*Q* filters suffer from degraded inband insertion loss (IL), passband flatness, and frequency selectivity, especially for narrowband applications.

In some applications, small inband magnitude variation, i.e., flat passband, is highly required while the additional IL can be tolerated, such as in a satellite input multiplexer and alternative receiver architectures [1]. In order to achieve flat passband and good selectivity with low-*Q* structures, the synthesis methods of bandpass filters with the consideration of loss have drawn much attention. Some techniques are proposed. The predistortion technique achieves good passband flatness by reflecting more power in the middle passband [2], resulting significantly deteriorated return loss (RL). The nonuniform-*Q* filters have improved the inband performance by controlling the *Q*-factor ratio of the resonators [3], with the degradation of RL and transmission zeros. The lossy coupling filters have the same performances as the ideal ones except for the additional attenuation [4], [5], which may be compensated by preamplifiers. Thus, the lossy coupling method is one attractive candidate for the high-performance miniaturized filter design. For a symmetric response, the lossy filter can be designed only with purely reactive and resistive couplings [6], [7]. As synthesized in [5], complex couplings with both nonzero real and imaginary parts will be required when an asymmetric response with transmission zero(s) is desired. Thus, the com-

Manuscript received November 22, 2017; revised February 12, 2018; accepted April 15, 2018. Date of publication May 9, 2018; date of current version June 5, 2018. This work was supported by the National Science Foundations of China under Grant 61370008 and Grant 61327803. *(Corresponding author: Lin-Sheng Wu.)*

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Digital Object Identifier 10.1109/LMWC.2018.2829349



Fig. 1. Proposed transmission lines loaded with (a) series and (b) shunt resistors, and (c) equivalent model.

plex coupling is to be realized as the key issue to build up this kind of lossy filter.

In this paper, a substrate integrated waveguide (SIW) lossy filter is proposed to achieve an asymmetric response. Different from the Chebyshev SIW filter in [6] and [7], complex coupling coefficients should be applied in this paper, which are realized with transmission lines loaded with series and shunt resistors. The complex-coupled filter with one transmission zero is built up by uniform-*Q* resonators. Good performance is achieved.

#### II. ANALYSIS AND DESIGN

Fig. 1(a) and (b) shows the two-port circuit models for the complex couplings whose imaginary parts are positive and negative, respectively. Both circuits can be equivalent to a *J* -inverter with a complex immittance and two residual shunt complex admittances at two ports, as shown in Fig. 1(c).

The ABCD matrixes  $A_a$  of the circuit in Fig. 1(a) and  $A_c$ of its equivalent circuit in Fig. 1(c) are calculated by

$$
\mathbf{A}_a = \begin{bmatrix} \cos\theta + j\frac{R_s \sin\theta}{2Z_0} & \frac{1+\cos\theta}{2}R_s + jZ_0 \sin\theta\\ \frac{\cos\theta - 1}{2Z_0^2}R_s + j\frac{\sin\theta}{Z_0} & \cos\theta + j\frac{R_s \sin\theta}{2Z_0} \end{bmatrix} (1)
$$

$$
\mathbf{A}_c = \begin{bmatrix} jY_a/J_a & j/J_a \\ j(Y_a^2/J_a + J_a) & jY_a/J_a \end{bmatrix}.
$$
 (2)

By making (1) and (2) equal to each other, the equivalent admittance  $J_a$  of the  $J$ -inverter and the residual admittance  $Y_a$  of the shunt circuit shown in Fig. 1(c) for Fig. 1(a) are obtained by

$$
J_a = \frac{4Z_0 \sin \theta}{D_a} + j \frac{2R_s(1 + \cos \theta)}{D_a}
$$
 (3)

$$
Y_a = \frac{2R_s(1+\cos\theta)}{D_a} + j\frac{R_s^2\sin\theta(1+\cos\theta) - 2Z_0^2\sin 2\theta}{Z_0D_a}
$$
 (4)

where

$$
D_a = R_s^2 (1 + \cos \theta)^2 + 4Z_0^2 \sin^2 \theta.
$$
 (5)

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Similarly, the equivalent admittance  $J_b$  of the  $J$ -inverter in Fig. 1(b) and the residual admittance  $Y_b$  of the shunt circuit shown in Fig. 1(c) for Fig. 1(b) are deduced by

$$
J_b = \frac{4R_p^2 \sin \theta}{Z_0 D_b} - j \frac{2R_p (1 - \cos \theta)}{D_b}
$$
 (6)

$$
Y_b = \frac{2R_p(1 - \cos\theta)}{D_b} - j\frac{Z_0^2 \sin\theta (1 - \cos\theta) + 2R_p^2 \sin 2\theta}{Z_0 D_b} \tag{7}
$$

$$
D_b = Z_0^2 (1 - \cos \theta)^2 + 4R_p^2 \sin^2 \theta.
$$
 (8)

It is seen from (3) and (6) that the imaginary parts of inverter admittances for the circuits of Fig. 1(a) and (b) are positive and negative, respectively. Therefore, a flat-passband lossy filter with an asymmetric response can be realized by the circuits shown in Fig. 1(a) and (b) for the complex coupling coefficients.

The residual susceptance slope parameters  $b_a$  and  $b_b$  of the circuits shown in Fig.  $1(a)$  and  $(b)$  can be derived by calculating the derivatives of  $\text{Im}(Y_a)$  in (4) and  $\text{Im}(Y_b)$  in (7) as

$$
b_a = \frac{2Z_0^4 \sin^2 \theta + R_s^4 \cos^6 \frac{\theta}{2} - 2R_s^2 Z_0^2 \cos \theta \cos^4 \frac{\theta}{2}}{4Z_0 \left(R_s^2 \cos^4 \frac{\theta}{2} + Z_0^2 \sin^2 \theta\right)^2 / \theta} \tag{9}
$$

$$
b_b = \frac{2R_p^4 \sin^2 \theta + Z_0^4 \sin^6 \frac{\theta}{2} + 2R_p^2 Z_0^2 \cos \theta \sin^4 \frac{\theta}{2}}{4Z_0 \left(Z_0^2 \sin^4 \frac{\theta}{2} + R_p^2 \sin^2 \theta\right)^2 / \theta}.
$$
 (10)

The residual conductance and susceptance can be absorbed by the two adjacent resonators connected to the *J* -inverter. Thus, the resonant frequencies and susceptance slope parameters of the two resonators should be adjusted, accordingly, during the filter design. The resonant frequency *fr* and the susceptance slope parameter  $b_{\text{in}}$  of the resonator are modified by

$$
f_r \Leftarrow f_r[1 + 0.5\mathrm{Im}(Y_i)/b_{\mathrm{in}}] \tag{11}
$$

$$
b_{\rm in} \Leftarrow b_{\rm in} - b_i \tag{12}
$$

where the subscript *i* represents *a* or *b*.

### III. DESIGN EQUATIONS

A third-order generalized Chebyshev lossy filter with a flat passband is designed as an example, which is centered at  $f_0 = 5$  GHz with one transmission zero located at 5.39 GHz. The fractional bandwidth is  $FBW = 5%$  and the inband RL is 20 dB, and all the resonators share the same unloaded *Q*-factor of 450. The coupling matrix **M** of the filter is synthesized as [5], where NS and NL denote the nonresonating nodes connected to source and load nodes, respectively, and the normalized coupling coefficient between the nodes *j* and  $k$  is denoted by  $M_{jk}$ . The scaling factor, i.e. the coupling coefficient between the source and NS nodes (the load and NL nodes) [4], is  $M_{S,NS} = M_{NL,L} = 0.1898$ . The coupled resonator circuit model and its coupling scheme of the developed filter are shown in Fig. 2.

The complex coupling *M*<sup>13</sup> between the resonators 1 and 3 is realized with the structure shown in Fig. 1(a) while  $M_{12}$  and  $M_{23}$  are realized with the structure shown in Fig. 1(b), since the imaginary parts of  $M_{13}$  and  $M_{12}$ (*M*23) are positive and negative, respectively. The inverter admittances  $J_{13}$  and  $J_{12}$  ( $J_{23}$ ) are given as

$$
J_{13} = M_{13} \text{FBW} \sqrt{b_1^{(3)} b_3^{(1)}}
$$
 (13)

$$
J_{12} = M_{12} \text{FBW} \sqrt{b_1^{(2)} b_2^{(1)}} = J_{23} = M_{23} \text{FBW} \sqrt{b_2^{(3)} b_3^{(2)}} \quad (14)
$$

where  $b_j^{(k)}$  is the susceptance slope parameter  $b_{\text{in}}$  of the node *j* seen from the node *k*. By comparing (3) with (13), the loaded resistance  $R_{13}$  and the characteristic impedance  $Z_{13}$  of the microstrip line between the resonators 1 and 3 are obtained by

$$
R_{13} = 2\text{Im}(M_{13}) / \left[\text{FBW}\sqrt{b_1^{(3)}b_3^{(1)}}|M_{13}|^2(1+\cos\theta_{13})\right]
$$
 (15)  

$$
Z_{13} = \text{Re}(M_{13}) / \left(\text{FBW}\sqrt{b_1^{(3)}b_3^{(1)}}|M_{13}|^2\sin\theta_{13}\right)
$$
 (16)

where  $\theta_{ik}$  is the electrical length of microstrip line between nodes *j* and *k*. Similarly,  $R_{12}$  ( $R_{23}$ ) and  $Z_{12}$  ( $Z_{23}$ ) are

$$
R_{12} = R_{23} = \frac{(1 - \cos \theta_{12}) \text{Re}^2(M_{12})}{-2 \text{FBW} \sqrt{b_1^{(2)} b_2^{(1)}} |M_{12}|^2 \sin^2 \theta_{12} \text{Im}(M_{12})}
$$
(17)

$$
Z_{12} = Z_{23} = \text{Re}(M_{12}) / (\text{FBW} \sqrt{b_1^{(2)} b_2^{(1)}} |M_{12}|^2 \sin \theta_{12}). \quad (18)
$$

The characteristic impedance of the 90° lines between the source (load) and the nonresonating nodes is obtained by

$$
Z_{S,NS} = Z_{\text{NL},L} = 1/J_{S,NS} = \sqrt{Z_p} / (M_{S,NS} \sqrt{b_2^{\text{(NS)}} \text{FBW}})
$$
 (19)  
where  $Z_p$  is the reference impedance of the I/O ports. When

 $Z_{S,NS} = Z_{NL,L} = Z_p$ , the two line sections can be absorbed by ports and the reference planes of source and load are shifted to the locations of the NS and NL nodes, respectively. The characteristic impedance of the microstrip line between resonator 1 (3) and the nonresonating node NS (NL) is

$$
Z_{\rm NS,1} = Z_{3,\rm NL} = 1/J_{\rm NS,1} = 1/(M_{\rm NS,1} \rm FBW \sqrt{b_1^{\rm (NS)} b_2^{\rm (NS)}}). \tag{20}
$$

The resistance between resonator 2 and the nonresonating node NS (NL) is calculated by

$$
R_{\rm NS,2} = R_{2,\rm NL} = 1/[b_2^{\rm (NS)} FBW \text{Im}(-M_{\rm NS,2})].
$$
 (21)

The electrical length of the two sections of microstrip lines connected to the resistors  $R_{\text{NS,2}}$  and  $R_{\text{2,NL}}$  are 180° and 360°.

# IV. RESULTS AND DISCUSSION

Here, we choose  $b_1^{(3)} = b_3^{(1)} = b_2^{(2)} = b_2^{(1)} = b_2^{(3)} = b_3^{(3)} = b_3^{(2)} = b_3^{(3)} = b_3^{$ 0.35 S, and  $\theta_{12} = \theta_{23} = \theta_{13} = 90^\circ$  for simplification. The values of  $R_{13} = 39 \Omega$ ,  $Z_{13} = 154 \Omega$ ,  $R_{12} = R_{23} = 1500 \Omega$ , and  $Z_{12} = Z_{23} = 58 \Omega$  are obtained from (15)–(18). Set  $b_2^{(NS)} =$  $b_2^{(NL)} = 11.1$  S, then we have  $R_{\rm NS,2} = R_{2,\rm NL} = 383 \Omega$ 



derived by



Fig. 2. Coupled resonator circuit model of the developed filter, where the coupling between the nonresonating node NS (NL) and resonator 2 is purely resistive, the coupling between resonators 1 and 2 (2 and 3) is with a complex coefficient whose imaginary part is negative, related to Fig. 1(b), and the coupling between resonators 1 and 3 is with a complex coefficient whose imaginary part is positive, related to Fig. 1(a).



Fig. 3. Layout of the fabricated lossy filter with complex couplings. Geometrical parameters are  $L_0 = 16.24, L_1 = 13.64, L_2 = 7.9,$  $L_3 = 19.64, L_4 = 6.48, L_5 = 10.87, L_6 = 6.5, L_7 = 2, L_8 = 2, L_9 = 2.8,$  $L_{10} = 1.6, L_{11} = 2.6, L_{12} = 7.34, L_{13} = 10.12, L_{14} = 4.1, L_{15} = 8.62,$  $L_{C0}$  = 3.76,  $L_{C1}$  = 4.4,  $L_{C2}$  = 3.9,  $L_{C3}$  = 3.86,  $L_R$  = 3.6,  $W_R = 1.2, W_1 = 3.38, W_2 = 0.2, W_3 = 0.22, W_4 = 2.45, g_C = 0.4,$ and  $g_R = 1.2$  (unit: mm). The resistors (package size of 0603) are with  $R_{\text{NS},2} = R_{2,\text{NL}} = 390 \Omega$ ,  $R_{12} = R_{23} = 1500 \Omega$ , and  $R_{13} = 39 \Omega$ . The structure within the red solid circle is related to Fig. 1(a) while the structures within the blue dotted circles are related to Fig. 1(b).

from (21) and  $Z_{S,NS} = Z_{NL,L} = 50 \Omega$  from (19) to remove these two line sections. Let  $b_1^{(NS)} = b_3^{(NL)} = 0.4$  S, then  $Z_{\text{NS},1} = Z_{3,\text{NL}} = 47 \Omega$  is obtained from (20). According to (12), the susceptance slope parameters are slightly modified as  $b_1^{(3)} = b_3^{(1)} = 0.345$  S,  $b_1^{(2)} = b_2^{(1)} = b_2^{(3)} = b_3^{(2)} = 0.336$  S. The resonant frequencies of resonators 1–3 are 4.989, 5.044, and 4.989 GHz, respectively, according to (11). As shown in Fig. 3, the prototype is fabricated on a 40-mil-thick taconic TLY-5A substrate, with  $\varepsilon_r = 2.17$  and tan $\delta = 0.0009$ , where the unloaded *Q-*factor of the hexagonal SIW resonators is about 450, extracted from the full-wave electromagnetic simulation. The physical dimensions and resistances of the filter are shown in Fig. 3. The 0603 resistors  $R_{12}$ ,  $R_{23}$ , and  $R_{13}$ are mounted over slots on the top metal plane and the resistor  $R_{12}$  ( $R_{23}$ ) is grounded by the metallic vias between SIW cavities. Note that the 390- $\Omega$  resistors, which are available in practice, are used instead of the  $383-\Omega$  resistors.

The synthesized, simulated, and measured S-parameters of the prototype are plotted in Fig. 4. Good agreement is achieved among them. The measured response is centered at 4.97 GHz with one transmission zero located at 5.35 GHz. The measured frequency shift of 30 MHz from the synthesized and simulated ones is mainly due to the tolerance of permittivity. By changing the permittivity to  $\varepsilon_r = 2.2$  in simulation, the modified simulated response is also plotted in Fig. 4, which agrees well with the measured one. The difference of RLs between the measured and simulated results is mainly due to the fabrication tolerance. The measured minimum inband IL is 2.2 dB, which is almost the same as the synthesized and simulated values. The 0.2-dB bandwidth is about 279 MHz



Fig. 4. Synthesized (Syn.), simulated (Sim.), measured (Mea.), and modified simulated (Mod.) results of the filter prototype.

in measurement, which matches the synthesized and simulated ones well. This 0.2-dB bandwidth is equivalent to a lossless synthesized filter but realized with the uniform unloaded *Q*-factor of 1900. Compared with the lossless synthesized filter but realized with the uniform unloaded *Q*-factor of 450, the 0.2-dB bandwidth of the proposed lossy filter is broadened by 20%.

## V. CONCLUSION

A lossy SIW filter with an asymmetric response and complex couplings is developed to improve passband flatness with uniform-*Q* resonators. The method is established with detailed design equations to realize complex coupling coefficients with positive and negative imaginary parts based on resistorloaded transmission lines. An equivalent circuit model of the lossy filter is proposed and all the critical parameters can be obtained analytically. Good agreement is obtained among the synthesized, simulated, and measured results, which validates our idea.

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