# Compact Tunable Bandpass Filter With a Fixed Out-of-Band Rejection Based on Hilbert Fractals

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Abstract—This paper proposes a new type of compact tunable bandpass filter with bandwidth tuning and out-of-band fixed rejection. Because we employ the modified Hilbert fractal structure loaded with varactors as resonators, the tunable filter has a very compact configuration and a constant shape over the entire tuning range. The frequency selectivity is improved by introducing a cross coupling between the source and the load. As a result of the utilization of a pair of properly designed feedlines, the frequency tuning and the out-of-band rejection of the filter are independent of each other, which simplifies its operation significantly. Two filter prototypes have been realized with the same size of  $25.0 \times 17.0 \times 1.0 \text{ mm}^3$ . Their superior performances have been demonstrated experimentally, with good agreement obtained between their simulated and measured *S*-parameters.

*Index Terms*—Bandpass filter (BPF), fractal, Hilbert resonator, out-of-band rejection, tunable filter.

#### I. INTRODUCTION

**P**LANAR varactor-tunable filters with compact structures are very important for the development of multifunctional multiband RF and microwave systems. In the past two decades, their studies have been mainly focused on the frequency tuning function realized by resonant structures [1]–[5], such as combined-line and interdigital resonators. In order to further improve their performances for some integrated compact multiband radio devices, such as home eNodeB [6], both frequency and bandwidth tuning capabilities are required [7]–[18]. By using a stepped-impedance resonator (SIR), a varactor-tunable bandpass filter (BPF) has been realized with a volume of  $38 \times 38 \times 0.8$  mm<sup>3</sup> in [7]. Its frequency tuning range is 12.5% at 2.0 GHz, the variation of its 3-dB fractional bandwidth (FBW) is within 3.2%, and

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the insertion loss is 4–6 dB. Another type of planar varactortunable BPF, based on a set of independent electric and magnetic coupling structures, has been presented in [15], in which the filters have been designed with three different FBW variations. In [16], a planar varactor-tunable BPF has been developed with a novel dual-mode resonator, whose its frequency tuning range is 41% from 0.57 to 0.98 GHz.

The above tunable BPFs have frequency and/or bandwidth tuning capabilities. However, it is still very challenging to tune the frequency and bandwidth, suppress the fixed frequency, and miniaturize the volume simultaneously. These three features are very essential for compact advanced multimode and multiband radio devices, including wireless system-on-apackage (SoP) [8], with interference suppression, as indicated in [19] and [20]. In order to find out an appropriate planar structure for the tunable filters with high performance, we resort to the unique Hilbert fractal structure. It has been utilized to build up planar passives [21]-[24], such as antennas, resonators, and high-impedance surfaces, because of its very compact geometry, multiband resonance, and self-similar response. However, to the best of our knowledge, no research has been reported for using the Hilbert fractal structure to design varactor-tunable filters. That is the main motivation of this paper.

In this paper, one type of compact varactor-tunable BPF is proposed using the third-order Hilbert fractal (H<sub>3</sub>) resonator. Its frequency tuning is achieved by the loaded varactors, and its out-of-band rejection is accomplished by the external feedlines next to the H<sub>3</sub> resonators. As tuning and rejection are independent of each other, it leads to a simple scheme for both controls, which is very important for multimode and multiband wireless SoP system development. In our design, the H<sub>3</sub> resonator, together with loaded varactors, is carefully analyzed, according to the coupled-line model. Two prototypes have been fabricated and measured to verify our design.

The organization of the rest of this paper is as follows. Section II presents modeling and design of the  $H_3$  resonator and the planar varactor-tunable BPF, and the resonant characteristics of the tunable  $H_3$  resonator are examined numerically. In Section III, the measured and simulated performances are given and analyzed for the two fabricated filter prototypes. Conclusions are drawn in Section IV.

#### II. MODELING AND DESIGN

Fig. 1(a) shows the layout of the proposed tunable BPF, which consists of I/O coupled line sections and two modified

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Fig. 1. Proposed tunable filter. (a) Layout. (b) Coupling scheme.

 $H_3$  resonators. Each of the resonators is loaded by a varactor  $C_V$ . Points indicted by  $P_2$  and  $P_4$  are shorted or opened here.

The coupling scheme of the tunable BPF is shown in Fig. 1(b), where S and L represent its source and load nodes, respectively.

To model and design the filter, both even- and odd-mode input admittances  $Y_{F1}^e$  and  $Y_{F1}^o$ , seen from the TT' plane in Fig. 1(a), is formulated with a magnetic or an electric wall placed at the OO' plane, respectively

$$Y_{F1}^{e,o} = y_{p2} \frac{A^{e,o} y_{p1}^{e,o} + j B^{e,o} y_{p2}^{tap} \varphi_{p2}}{B^{e,o} y_{p2} + j A^{e,o} y_{p1}^{e,o} \tan \varphi_{p2}}$$
(1)

and

$$A^{e,o} = y_{F0}^{e,o} + j y_{p1}^{e,o} \tan \varphi_{p1}$$
(2)

$$B^{e,o} = y_{p1}^{e,o} + j y_{F0}^{e,o} \tan \varphi_{p1}$$
(3)

$$y_{F0}^{e,o} = \begin{cases} P_2 \text{ and } P_4 \text{ opened} \\ P_2 \text{ and } P_4 \text{ shorted} \end{cases}$$
(4)

where the superscripts *e* and *o* represent the even and odd modes, respectively. By looking into the H<sub>3</sub> resonator from the UU' plane in Fig. 1(a), the input admittances  $y_o^{e,o}$  and  $y_s^{e,o}$  can be extracted with electromagnetic (EM) simulation. Further, using the circuits approach, the input admittances  $Y_F^e$ and  $Y_F^o$  at the input port can be obtained from  $Y_{F1}^e$  and  $Y_{F1}^o$ , respectively. The overall admittance matrix of the above filter can then be described by

$$\begin{bmatrix} y_{F11} & y_{F12} \\ y_{F21} & y_{F22} \end{bmatrix} = \frac{1}{2} \begin{bmatrix} Y_F^e + Y_F^o & Y_F^e - Y_F^o \\ Y_F^e - Y_F^o & Y_F^e + Y_F^o \end{bmatrix}.$$
 (5)

In our design, three important issues have to be taken care of, which are as follows.

- 1) In order to provide a nearly constant filter response shape and bandwidth, the separation between the oddand even-mode resonant frequencies  $f_0^o$  and  $f_0^e$  of the H<sub>3</sub> resonator should be kept almost unchanged when changing the capacitance  $C_V$ .
- 2) The derivative of the input susceptance to  $f_0^o$  and  $f_0^e$  for H<sub>3</sub> resonator's external coupling should be nearly constant so that the filter has a constant bandwidth over the entire frequency tuning range.
- 3) Rejection at a fixed frequency is realized by the fixed transmission zero at the frequency  $f_z$ . Therefore, the input admittances  $Y_F^e = Y_F^o$  at  $f_z$  is required, as the capacitance  $C_V$  changes.

### A. H<sub>3</sub> Resonator

With an electric or a magnetic wall inserted at the OO' plane in Fig. 1(a), two circuits for the odd and even modes with no external coupling are obtained. Their layouts and approximate equivalent circuit models are shown in Fig. 2(a)–(c). For simplicity, we assume that the H<sub>3</sub> resonator consists of transmission lines (TLs), in which the *i*th TL has electrical length of  $\theta_i$ . The parasitic and coupling effects of the *i*th TL within the resonator are considered by reducing the  $\theta_{0i}$  with  $\Delta \theta$ , i.e.,  $\theta_i = \theta_{0i} - \Delta \theta$ , where  $\Delta \theta$  is about 3.1° in our case and  $\theta_{0i}$  (i = 1, 2, 3, 4) is the electrical length of the *i*th TL without parasitic and coupling effects.

In Fig. 2, we add the reference port  $P_1$  to calculate the oddand even-mode input admittances of the  $H_3$  resonator

$$Y_{in} = \frac{a + (Y_0 \tan \theta_3) b + [(Z_0 \tan \theta_3) a + b] Y_l}{c + (Y_0 \tan \theta_3) d + [(Z_0 \tan \theta_3) c + d] Y_l}$$
(6)

$$Y_l = j\omega C_V + \begin{cases} +jY_0 \tan \theta_4, P_3 \text{ and } P_4 \text{ opened} \\ -jY_0 \cot \theta_4, P_3 \text{ and } P_4 \text{ shorted} \end{cases}$$
(7)

$$a = C_1 A_2 + D_1 C_2 \tag{8}$$

$$b = C_1 B_2 + D_1 D_2 \tag{9}$$

$$= A_1 A_2 + B_1 C_2 \tag{10}$$

$$d = A_1 B_2 + B_1 D_2 \tag{11}$$

where  $Z_0 = 1/Y_0$ , with  $Y_0$  being the characteristic admittances of the transmission lines for the H<sub>3</sub> resonator

$$\begin{bmatrix} A_i & B_i \\ C_i & D_i \end{bmatrix} = \begin{bmatrix} 1 + Z_0 Y \tan \theta_i \tan \theta & Z_0 \tan \theta_i + Z \tan \theta \\ Y_0 \tan \theta_i + Y \tan \theta & \tan \theta (Y_0 Z \tan \theta_i + 1) \end{bmatrix}.$$
(12)

Equations (6)–(12) can be used to describe both odd- and even-mode characteristics. For the odd mode, we substitute  $\theta = \theta^o$  and  $Y = Y^o$  into (12), while for the even mode, we substitute  $\theta = \theta^e$  and  $Y = Y^e$ .  $Y^o$  and  $Y^e$  represent the odd- and even-mode characteristic admittances of the coupled transmission lines, respectively.  $\theta^o$  and  $\theta^e$  are their



Fig. 2. (a) Layout of the  $H_3$  resonator with an electric and a magnetic wall inserted at the OO' plane. (b) Its odd- and (c) even-mode equivalent circuits.

corresponding electrical lengths. When  $C_V = 0$ , the initial length of the single H<sub>3</sub> resonator is determined by its central frequency  $f_0$ .

At  $f_0^o$  and  $f_0^e$ , the following resonant conditions should be satisfied:

$$Im[Y_{ino}(f_0^o)] = 0 (13)$$

$$Im[Y_{ine}(f_0^e)] = 0.$$
 (14)

Fig. 3 shows the effectiveness of the derived equations to deal with parasitic and coupling effects in the even and odd mode and the variations of the odd- and even-mode resonant frequencies for different values of  $C_V$ , with given parameters. It is observed that, for the designed two H<sub>3</sub> resonators, when  $C_V$  varies from 0 to 6 [pF], the separation between  $f_0^o$  and  $f_0^e$  is almost unchanged.

The coupling coefficient k between two H<sub>3</sub> resonators is calculated by

$$k = \frac{\left| (f_0^e)^2 - (f_0^o)^2 \right|}{\left| (f_0^e)^2 + (f_0^o)^2 \right|} \approx \frac{\left| f_0^e - f_0^o \right|}{f_0}.$$
 (15)

Fig. 4 shows the coupling coefficient k as a function of the spacing  $W_m$  between two H<sub>3</sub> resonators when P<sub>3</sub> and P<sub>4</sub> are opened and  $C_V = 0$ . This result is extracted from the full-wave EM simulation, and the analytical one is also provided for comparison.

The relative tuning range is defined as the ratio of the frequency tuning range  $(f_2 - f_1)$  to the capacitance variation  $(C_{V2} - C_{V1})$  when the loading capacitance  $C_V$  varies from



Fig. 3. Resonant frequencies of the odd and even modes as a function of varactor capacitance  $C_V$  when P<sub>2</sub> is opened ( $W_m = 0.15$ ,  $W_0 = W_1 = 0.5$ ,  $W_2 = 1.02$ ,  $L_1 = 2.05$ , and  $L_2 = 0.52$  mm).



Fig. 4. Coupling coefficient k as a function of the spacing  $W_m$  between two H<sub>3</sub> resonators when P<sub>2</sub> and P<sub>4</sub> are opened and  $C_V = 0$  pF ( $W_m = 0.15$ ,  $W_0 = W_1 = 0.5$ ,  $W_2 = 1.02$ ,  $L_1 = 2.05$ , and  $L_2 = 0.52$  mm).

 $C_{V1} = 0.5$  pF to  $C_{V2} = 5.5$  pF. Here,  $f_1$  and  $f_2$  are the lowest and highest central frequencies of the tunable H<sub>3</sub> resonator, respectively. To demonstrate how to control the tuning range, the dependence of the tuning frequency on the characteristic admittances  $Y_0$  and electrical lengths of H<sub>3</sub> resonators are investigated. As the odd-mode electrical length  $\theta^\circ = 140^\circ$ at 1.9 GHz, the tuning range with  $Y_0$  is calculated for the odd-mode case, which is plotted in Fig. 5(a). In addition, the variation of the tuning range with different  $\theta^\circ$  at 1.9 GHz and  $Y_0 = 0.009$  is plotted in Fig. 5(b). Based on Fig. 5, the proper linewidth  $W_1$  and the total length of H<sub>3</sub> resonator are obtained for the maximum relative tuning range.

# B. External Coupling of the H<sub>3</sub> Resonator

Fig. 6 shows the H<sub>3</sub> resonator together with its external coupling structure. The frequency-dependent susceptance slope parameter of  $Y_{F1}^{e}$  and  $Y_{F1}^{o}$  for this structure can be computed from the definition of external quality factor  $Q_{exe}$ . For the odd



Fig. 5. Variation of the relative tuning range of the odd mode with (a)  $Y_0$  and (b)  $\theta^{\circ}$ , where  $C_V$  is varied from 0.5 to 5.5 pF and P<sub>2</sub> shorted ( $L_{p0} = 12.9$ ,  $W_{p0} = 0.23$ , and  $S_{p0} = 0.25$  mm).

mode, we have

$$Q_{ex}^{o} = \frac{f_0^{o}}{\Delta f_0^{o}} = \frac{Z_{\text{port}} f_0^{o}}{2} \frac{\partial}{\partial f} \text{Im} \left[ Y_{F1}^{o} \left( f_0^{o} \right) \right]$$
(16)

where  $Z_{\text{port}} = 50 \ \Omega$  is the port impedance, and  $Q_{ex}^o$  and  $\Delta f_0^o$  are the external *Q*-factor and the 3-dB bandwidth of the odd mode, respectively. It can be obtained by

$$\frac{\partial}{\partial f} \operatorname{Im} \left[ Y_{F1}^{o} \left( f_{0}^{o} \right) \right] = \frac{2}{Z_{\text{port}} \Delta f_{0}^{o}}.$$
 (17)

Similarly, for the even mode, we have

$$\frac{\partial}{\partial f} \operatorname{Im} \left[ Y_{F1}^{e} \left( f_{0}^{e} \right) \right] = \frac{2}{Z_{\text{port}} \Delta f_{0}^{e}}$$
(18)

where  $\Delta f_0^e$  is the 3-dB bandwidth of the even mode.

It can be seen from (17) and (18) that  $\Delta f_0^o$  and  $\Delta f_0^e$  are unchanged only when the derivative of the input susceptance to  $f_0^o$  and  $f_0^e$  keeps constant over the entire tuning range. In other words, Im $[Y_{F1}^e]$  and Im $[Y_{F1}^o]$  should be the linear functions of  $f_0^o$  and  $f_0^e$  when the loading capacitance  $C_V$  varies from 0.5 to 5.5 pF.



Fig. 6. H<sub>3</sub> resonator together with the external coupling structure.



Fig. 7. Derivatives of the input susceptance to the resonated frequency when  $C_V$  varies from 0.5 to 5.5 pF and P<sub>2</sub> shorted ( $W_m = W_{mp1} = 0.15$ ,  $W_0 = W_1 = 0.5$ ,  $W_2 = 1.02$ ,  $L_1 = 2.05$ ,  $L_2 = 0.52$ ,  $L_{p0} = 12.9$ ,  $W_{p0} = 0.23$ , and  $S_{p0} = S_{p1} = 0.25$  mm).

In order to satisfy the conditions (17) and (18) for all values of  $C_V$  at the operated tuning range, appropriate initial values of  $L_{p0}$ ,  $S_{p0}$ , and  $W_{p0}$  are required. This can be achieved by properly choosing the geometrical parameters  $W_{p0}$ ,  $S_{p0}$ , and  $S_{p1}$  based on EM simulations. In general, for a given  $L_{p0}$ , which mainly depends on the transmission zero, a small value of  $S_{p0}$  or  $S_{p1}$  can provide strong coupling, and the external coupling can be enhanced with  $W_{p0}$  decreasing. Fig. 7 gives the desired results obtained from (17) and (18), and the extracted one from the EM simulations.

#### C. Source-Load Cross Coupling

To improve the filter frequency selectivity, it is important to properly introduce and design the source-load cross coupling.



Fig. 8. Source-load cross coupling of the filter. (a) Original circuit model. (b) Simplified circuit model.

Such a design is provided through a section of coupled lines, denoted by  $L_{p1}$ ,  $W_{p1}$ , and  $W_{mp1}$  in Fig. 1(a). Because of the small value of  $L_{p1}$ , the coupled lines can be approximated by a pair of noncoupled lines, of length  $L_{p1}$  and width  $W_{p1}$ , with a coupling capacitance  $C_c$  in parallel connection, as shown in Fig. 8.

In Fig. 8(b),  $C_c$  is extracted by using the parallel network theory, and

$$Y_{F2,c} = \begin{bmatrix} 1 & 1/j\omega C_c \\ 0 & 1 \end{bmatrix} + Y_{F2,n}$$
(19)

where

$$Y_{F2,c} = \frac{1}{2} \begin{bmatrix} Y_{F2,c}^{e} + Y_{F2,c}^{o} & Y_{F2,c}^{e} - Y_{F2,c}^{o} \\ Y_{F2,c}^{e} - Y_{F2,c}^{o} & Y_{F2,c}^{e} + Y_{F2,c}^{o} \end{bmatrix}$$
(20)

$$Y_{F2,n} = \frac{1}{2} \begin{bmatrix} Y_{F2,n}^{e} + Y_{F2,n}^{o} & Y_{F2,n}^{e} - Y_{F2,n}^{o} \\ Y_{F2,n}^{e} - Y_{F2,n}^{o} & Y_{F2,n}^{e} + Y_{F2,n}^{o} \end{bmatrix}.$$
 (21)

For the coupled lines  $L_{p1}$ , the input admittances of the even and odd modes  $Y_{F2,c}^{e,o}$  are calculated by

$$Y_{F2,c}^{e,o} = y_p \frac{y_{F0}^e + j y_{p1}^{e,o} \tan \phi_1}{y_{p1}^{e,o} + j y_{F0}^e \tan \phi_1}.$$
 (22)

TABLE I DIMENSIONS OF THE TUNABLE FILTER (mm)

$W_0$	$W_1$	$W_2$	$L_1$	$L_2$	$W_m$	$W_{p0}$	$L_{p0}$	$S_{p0}$
0.5	0.5	1.02	2.05	0.52	0.15	0.23	12.9	0.25
$W_{p1}$	$W_{p2}$	$L_{p0}$	$L_{p1}$	$L_{p2}$	$L_{p3}$	$L_{p4}$	$S_{p1}$	$W_{p4}$
0.3	0.3	12.9	2.54	5.87	2.84	3.66	0.25	2.44
$W_m = W_{mp1}$ .								



Fig. 9. S-parameters of the filter obtained with EM simulations and simplified circuit model with P<sub>3</sub> and P<sub>4</sub> shorted ( $W_m = W_{mp1} = 0.15$ ,  $W_0 = W_1 = 0.5$ ,  $W_2 = 1.02$ ,  $L_1 = 2.05$ ,  $L_2 = 0.52$ ,  $L_{p0} = 12.9$ ,  $W_{p0} = 0.23$ , and  $S_{p0} = S_{p1} = 0.25$  mm).

Similarly, for a noncoupled line, the input admittance  $Y_{F2,n}^{e,o}$  is

$$Y_{F2,n}^{e,o} = y_p \frac{y_{F0}^{e,o} + jy_{p1} \tan \phi_1}{y_{p1} + jy_{F0}^{e,o} \tan \phi_1}.$$
 (23)

When the parameters shown in Table I and (19) are used, the value of  $C_c = 0.057$  pF is obtained. To validate our simplified model, full-wave simulations are carried out, with the results plotted in Fig. 9.

#### D. Fixed Out-of-Band Rejection

At the fixed out-of-band rejection frequency  $f_z$ , which is far from the tuning range of the resonator, the EM coupling between the feed line  $L_{p0}$  and the H<sub>3</sub> resonator is very weak, due to their frequency detuning. Therefore, for the given width,  $W_{p0}$  and  $W_{p1}$  of the microstrip line, the filter's response characteristics at  $f_z$  are mainly dominated by its end-opened line length  $L_p$ , made of  $L_p = L_{p0} + L_{p1}$  (see Fig. 6). In addition, this frequency detuning leads to the characteristics of H<sub>3</sub> resonator with loaded varactors to be almost independent of the filter performance at  $f_z$ . This means that the frequency tuning is independent of its out-of-band rejection. This feature of the filter can be used to design the tunable filter with out-of-band fixed rejections as in the following approaches.

Once  $\phi_{p1}$  (or  $L_{p1}$ ) and  $W_{p1}$  are determined by the needed source and load coupling first (frequency tuning), the  $\phi_{p0}$ 



Fig. 10. Central frequency  $f_0$  and rejection frequency  $f_z$  of the filter as functions of (a) capacitance  $C_V$  and (b) length  $L_{p0} = 7.92 + L_{p01}$  ( $W_m =$  $W_{mp1} = 0.15, W_0 = W_1 = 0.5, W_2 = 1.02, L_1 = 2.05, L_2 = 0.52, W_{p0} = 0.15, W_{p1} = 0$ 0.23, and  $S_{p0} = S_{p1} = 0.25$  mm).

(or  $L_{p0}$ ) can then be determined by using transmission zero condition of the end-opened lines  $L_p$ ,  $y_{in}^e(f_z) = y_{in}^o(f_z)$ without H<sub>3</sub> resonators existence

$$y_{p1}^{e}\tan(\phi_{p0}^{e}+\phi_{p1}) = y_{p1}^{o}\tan(\phi_{p0}^{o}+\phi_{p1})$$
(24)

and

$$\phi_{p0}^{e,o} = \tan^{-1} \left( \frac{y_{p0} \tan \phi_{p0}}{y_{p1}^{e,o}} \right)$$
(25)

where  $y_{p0}$ ,  $\phi_{p0}$ ,  $y_{p1}^e$ ,  $y_{p1}^o$ , and  $\phi_{p1}$  are as shown in Fig. 1(a). Corresponding to  $L_{p0} = 12.9$  mm and  $L_{p1} = 2.54$  mm when P<sub>2</sub> and P<sub>4</sub> ports are shorted, Fig. 10(a) shows the central frequency  $f_0$  and  $f_z$  against the capacitance  $C_V$  varying from 1.0 to 5.0 pF. Fig. 10(b) shows the variations of  $f_0$  and  $f_z$  for different values of  $L_{p01}$  [see Fig. 1(a)] when  $C_V = 2$  pF. The former shows that the  $f_z$  is nearly a constant for the given  $L_{p0}$  and  $L_{p1}$  over the tuning range. The latter shows that the longer the length  $L_{p01}$ , the smaller the  $f_z$  obtained.

Our design procedure of the tunable filter is summarized as follows.

TABLE II PARAMETERS FOR FILTER PROTOTYPES A AND B

Samples	А	В
Size (mm <sup>3</sup> )	$20.0\times13.0\times1.0$	$20.0 \times 13.0 \times 1.0$
Frequency tunability	13%	7.4%
Insertion loss (dB)	1.7	4.5
3-dB FBW	9.3%	11.3%
Central frequency (GHz)	1.71	1.95
Rejection (dB)	-35@2.4 GHz	-28@2.48 GHz





(b)

Fig. 11. Photographs of the filter prototypes. (a) Filter A. (b) Filter B.

- 1) According to the upper limit of the given frequency tuning range and the given minimum value of  $C_V$ , the initial parameters of two H<sub>3</sub> resonators ( $W_0$ ,  $W_1$ ,  $W_2$ ,  $L_1$ , and  $L_2$ ) are achieved.
- 2) As the BPF characteristics are given, the coupling coefficients  $k_{12}$  between two H<sub>3</sub> resonators, and  $k_{SL}$  between the source and the load, can be synthesized based on the theory of elliptical filters with double poles.
- 3) According to the desired value of  $k_{12}$ , the spacing  $W_m$ between the two H<sub>3</sub> resonators is determined.
- 4) Based on the desired value of  $k_{SL}$ , the parameters  $y_{n1}^{e,o}$ and  $\phi_{p1}$ , i.e.,  $L_{p1}$ ,  $W_{p1}$ , and  $W_{mp1}$ , can be determined. Furthermore, the geometrical parameters  $W_{p0}$ ,  $S_{p0}$ ,  $S_{p1}$ , and the initial  $L_{p0}$  are obtained by minimizing the derivatives of both even- and odd-mode input susceptances.
- 5) The length  $L_{p1}$  of the external lines is finally determined by the required out-of-band rejection at  $f_z$  over the entire tuning range.



Fig. 12. Simulated and measured *S*-parameters of filter A. (a) Reflection coefficients. (b) Transmission coefficients.

### **III. RESULTS AND DISCUSSION**

To validate our design, two filter prototypes, denoted A and B, are fabricated on a Taconic TSM-30 substrate, which has the relative permittivity of 3.0 and the thickness of 1.016 mm. Their dimensions and main performance parameters are listed in Tables I and II, respectively, and their photographs are shown in Fig. 11.

In their realizations, the RF circuits, including two H<sub>3</sub> resonators, two external coupled feedlines, and two loaded Infineon BB857 varactors [26], are located on the top PCB layer. The d.c. bias circuits shown in Fig. 1(a), consisting of two AVX chip capacitors  $C_0 = 15.4$  and  $C_d = 30$  pF, a resistor  $R_0 = 1.0 \text{ M}\Omega$ , and an AVX chip inductor  $L_{gd} = 39$  nH, are on its bottom layer. Here,  $C_0$  is series-connected with the varactor capacitance  $C_{V0}$ ,  $C_d$  is a decoupling capacitor,  $R_0$  isolates the d.c. from the RF signal, and  $L_{gd}$  is used as the RF choke. Four metallic via-holes, each with a diameter of 0.25 mm, are utilized to connect both RF and d.c. bias circuits.



Fig. 13. Measured insertion loss, 3-dB bandwidth, and central frequency of filter A with different bias voltages.



Fig. 14. Measured harmonic response of filter A.

# A. Open-Ended Tunable Filter A

As shown in Fig. 11(a), filter A has an overall volume of  $25.0 \times 17.0 \times 1.0 \text{ mm}^3$  with P<sub>2</sub> and P<sub>4</sub> open ended. Fig. 12 shows its simulated and measured S-parameters. With the bias voltage of 5 V, the central frequency is about 1.71 GHz, and the 3-dB FBW is 9.3%, with an in-band insertion loss of 1.7 dB, as shown in Fig. 13. When the bias voltage is increased to 25 V, the central frequency moves to 1.93 GHz, and the 3-dB FBW is 9.9%, with an insertion loss of 1.24 dB. In Fig. 13, the measurement shows an almost constant FBW, and the variation is within 0.6%, i.e., the frequency of filter A can be tuned with a nearly constant FBW. The measured inband return loss is always better than 10 dB over the entire tuning range. At both the lower and upper stop bands, the out-of-band rejections are better than 19 dB. The variation of insertion loss, when the central frequency moves from 1.71 to 1.93 GHz, is smaller than 0.8 dB.

Fig. 14 shows the measured harmonic response of filter A. It is observed that the bias voltage has little impact on its frequency responses in the lower stopband from 0.1 to 1.3 GHz



Fig. 15. Simulated and measured *S*-parameters of filter B. (a) Reflection coefficients. (b) Transmission coefficients.

and the upper one from 2.5 to 5.0 GHz. Their rejections are lower than -25.0 and -16.0 dB. There is a fixed transmission zero at 2.4 GHz, independent of the varactor capacitance. Such a property is useful for interference suppression.

#### B. Short-Ended Tunable Filter B

As shown in Fig. 11(b), filter B is has the same volume as filter A.  $P_2$  and  $P_4$  are short-ended in this case. Fig. 15 shows good agreement between its measured and simulated *S*-parameters. Its frequency tuning range is from 1.98 to 2.4 GHz; the measured in-band return and insertion losses are larger than 12 and 4.5 dB over the entire operating bandwidth, respectively. The insertion loss of filter B is larger than that of filters A, mainly due to the reduction of its FBW.

As shown in Fig. 16, when the bias voltage varies from 2 to 14 V, the central frequency of filter B is increased from 1.71 to 2.25 GHz, the 3-dB FBW is decreased from 7.6% to 5.7%, and the insertion loss is decreased from 5.7 to 3.5 dB. In other words, the FBW of filter B is decreased with its tuning frequency increasing.



Fig. 16. Measured insertion loss, 3-dB bandwidth, and central frequency of filter B for different bias voltages.



Fig. 17. Measured harmonic response of filter B.

Fig. 17 shows the measured harmonic response of filter B with the bias voltage varying from 7 to 14 V. A parasitic response arises around 1.1 GHz, with the level about -10 dB. It is because, at the frequency around 1 GHz, the H<sub>3</sub> resonator of filter B is operated in its nearly  $1/4\lambda_g$  resonant mode, which leads to the spurious passband. At around 2 GHz, the resonator is operated in its nearly  $3/4\lambda_g$  resonant mode.

The ratio of their resonant frequencies deviates from 3 because of the capacitive loading effect. Similar to filter A, there is also little impact of bias voltage on the frequency responses of filter B in the lower stop band from 0.1 to 1.0 GHz and the upper one from 2.5 to 5.0 GHz. The harmonic responses in the lower and upper stop bands are lower than -25.0 and -16.0 dB, respectively. At about 2.48 GHz, there is also a fixed transmission zero. The performances of filter A and B compared with the previously published ones are shown in Table III. It is observed that our tunable BPFs are very compact in size and also have a fixed out-of-band rejection. These features are very

	TABLE III
D	DETAILED COMPARISON OF THE PROPOSED FILTERS WITH PREVIOUSLY PUBLISHED FILTERS

Ref.	Frequency Tuning Range (GHz)	Bandwidth Tuning Range	IL (dB)	Fixed Out-of-Band Reject	<b>Normalized</b> Area $(\lambda_0^2)$	Feature
[2]	1.00-1.33	15%	$\leq$ 3.0	No	0.04	1
[7]	1.85-2.15	3.2%	$\leq 6.0$	No	0.17	1
[10]	2.10-2.70	0.3%	$\leq$ 5.0	No	0.67	2
[13]	0.47-0.86	33%	≤4.0	No	0.032	*
[15]	0.85-1.40	5.4%	≤2.9	No	0.036	2
	0.91-1.34	5.2-2.9%	$\leq 2.9$	No	0.036	1
	0.86-1.41	4.3-6.5%	≤3.5	No	0.036	3
[16]	0.60-1.07	14-8.2%	$\leq 1.8$	No	*	3
	0.57-0.98	16-9.2%	$\leq 2.2$	No	*	3
[17]	1.40-2.00	9%	$\leq 4.0$	No	0.033	*
Filter A	1.71-1.93	9.3-9.9%	≤1.7	-35 dB	0.021	2
Filter B	1.92-2.25	7.6–5.7%	≤4.1	-30 dB	0.021	4

1: Constant BW. 2: Constant FBW. 3: Increased FBW. 4: Decreased FWB. \* No data.

useful for multimode and multiband wireless SoP system applications.

# IV. CONCLUSION

In this paper, a new compact planar tunable filter was proposed with tunable frequency and fixed out-of-band rejection. The filter design is based on the Hilbert fractal structures loaded with varactors. The frequency tuning and its out-ofband rejection are independent of each other, which makes its operation simple. The circuit model and design procedure of the tunable filter were presented. Two filter prototypes were realized to validate our design. Filter A had a nearly constant FBW, while filter B had a decreasing FBW with increasing bias voltage of the varactor. Both of them had fixed out-of-band transmission zero, which could be utilized for interference suppression. Good agreement was obtained between the measured and simulated results for each filter prototype, and their frequency tuning and rejection capabilities were demonstrated.

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